

travel in both directions along the string, will be reflected at the ends, and will travel back in the opposite direction. Most of these waves interfere with each other and quickly die out. However, those waves that correspond to the resonant frequencies of the string will persist. The ends of the string, since they are fixed, will be nodes. There may be other nodes as well. Some of the possible resonant modes of vibration (standing waves) are shown in Fig. 11–40b. Generally, the motion will be a combination of these different resonant modes, but only those frequencies that correspond to a resonant frequency will be present.

To determine the resonant frequencies, we first note that the wavelengths of the standing waves bear a simple relationship to the length  $L$  of the string. The lowest frequency, called the **fundamental frequency**, corresponds to one antinode (or loop). And as can be seen in Fig. 11–40b, the whole length corresponds to one-half wavelength. Thus  $L = \frac{1}{2}\lambda_1$ , where  $\lambda_1$  stands for the wavelength of the fundamental frequency. The other natural frequencies are called **overtones**; for a vibrating string they are whole-number (integral) multiples of the fundamental, and then are also called **harmonics**, with the fundamental being referred to as the **first harmonic**.<sup>†</sup> The next mode of vibration after the fundamental has two loops and is called the **second harmonic** (or first overtone), Fig. 11–40b. The length of the string  $L$  at the second harmonic corresponds to one complete wavelength:  $L = \lambda_2$ . For the third and fourth harmonics,  $L = \frac{3}{2}\lambda_3$ , and  $L = 2\lambda_4$ , respectively, and so on. In general, we can write

$$L = \frac{n\lambda_n}{2}, \quad \text{where } n = 1, 2, 3, \dots$$

The integer  $n$  labels the number of the harmonic:  $n = 1$  for the fundamental,  $n = 2$  for the second harmonic, and so on. We solve for  $\lambda_n$  and find

$$\lambda_n = \frac{2L}{n}, \quad n = 1, 2, 3, \dots \quad (11-19a)$$

To find the frequency  $f$  of each vibration we use Eq. 11–12,  $f = v/\lambda$ , and we see that

$$f_n = \frac{v}{\lambda_n} = n \frac{v}{2L} = nf_1, \quad n = 1, 2, 3, \dots, \quad (11-19b)$$

where  $f_1 = v/\lambda_1 = v/2L$  is the fundamental frequency. We see that each resonant frequency is an integer multiple of the fundamental frequency.

Because a standing wave is equivalent to two traveling waves moving in opposite directions, the concept of wave velocity still makes sense and is given by Eq. 11–13 in terms of the tension  $F_T$  in the string and its mass per unit length ( $m/L$ ). That is,  $v = \sqrt{F_T/(m/L)}$  for waves traveling in both directions.

**EXAMPLE 11–14 Piano string.** A piano string is 1.10 m long and has a mass of 9.00 g. (a) How much tension must the string be under if it is to vibrate at a fundamental frequency of 131 Hz? (b) What are the frequencies of the first four harmonics?

**APPROACH** To determine the tension, we need to find the wave speed using Eq. 11–12 ( $v = \lambda f$ ), and then use Eq. 11–13, solving it for  $F_T$ .

**SOLUTION** (a) The wavelength of the fundamental is  $\lambda = 2L = 2.20$  m (Eq. 11–19a with  $n = 1$ ). The speed of the wave on the string is  $v = \lambda f = (2.20 \text{ m})(131 \text{ s}^{-1}) = 288 \text{ m/s}$ . Then we have (Eq. 11–13)

$$F_T = \frac{m}{L} v^2 = \left( \frac{9.00 \times 10^{-3} \text{ kg}}{1.10 \text{ m}} \right) (288 \text{ m/s})^2 = 679 \text{ N}.$$

(b) The frequencies of the second, third, and fourth harmonics are two, three, and four times the fundamental frequency: 262, 393, and 524 Hz.

**NOTE** The speed of the wave on the string is *not* the same as the speed of the sound that is produced in the air (as we shall see in Chapter 12).

<sup>†</sup>The term “harmonic” comes from music, because such integral multiples of frequencies “harmonize.”

Fundamental frequency

Overtones and harmonics